# **RESEARCH ARTICLE**



# Aftershocks following the 2011 Tohoku-Oki earthquake driven by both stress transfer and afterslip



# Abstract

Aftershocks are a fundamental characteristic of seismicity, and their generation mechanism is mainly characterized by two physical models, stress transfer from large earthquakes and afterslip-induced stress loading. However, the contribution of each mechanism to aftershock generation remains unclear. Here we investigate the spatiotemporal variations in aftershock activity following the 2011 Tohoku-Oki earthquake by applying the Hierarchical Space–Time Epidemic-Type Aftershock Sequence (HIST-ETAS) model to the decade of recorded seismicity since the mainshock. Using the estimated HIST-ETAS model, we categorize the aftershocks into background earthquakes (which are caused by aseismic phenomena) and triggered earthquakes (which are caused by earthquake-to-earthquake interactions). Most of the earthquakes that occurred updip of the large coseismic slip zone along the Japan Trench are triggered earthquakes, consistent with the lack of afterslip in this area. Conversely, background earthquakes are the predominant earthquake type in the long-term downdip of the large coseismic slip zone, and they positively correlate with the afterslip evolution. Our results suggest the importance of combining these two end-member aftershock generation models to explain aftershock activity and thus provide new insights into the relationship between afterslip in a given region that hosts a large earthquake, particularly where geodetic observation networks are too sparse to evaluate afterslip evolution.

**Keywords** Aftershock, Stress transfer, Afterslip, 2011 Tohoku-Oki earthquake, HIST-ETAS model, Background earthquakes, Triggered earthquakes

## **1** Introduction

The 2011 M9 Tohoku-Oki earthquake occurred on March 11, 2011, along the boundary between the subducting Pacific Plate and the overriding Okhotsk Plate (Fig. 1). Previous studies reveal the aftershock distribution in and around the focal area of this megathrust earthquake (e.g., Asano et al. 2011; Kato and Igarashi 2012). Asano et al. (2011) estimated the centroid moment tensors of earthquakes before and after the 2011 Tohoku-Oki earthquake and revealed that interplate earthquakes did not occur within the large coseismic slip area of this megathrust earthquake, whereas both interplate and non-interplate earthquakes occurred downdip of the large coseismic slip area and along the Japan Trench, as shown in Fig. 1b. Asano et al. (2011) also suggested that the interplate earthquakes that occurred downdip of the coseismic slip area were primarily loaded by coseismic slip and would be promoted by postseismic slip, with normal fault-type



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**Fig. 1** Epicenter distributions of the detected earthquakes (01/10/1997–20/02/2021) in the study region. **a** Earthquake distribution ( $M_j \ge 4.0$ ,  $\le 100$  km depth) before the 2011 Tohoku-Oki earthquake (01/10/1997–11/03/2011). **b** Earthquake distribution ( $M_j \ge 4.0, \le 100$  km depth) after the 2011 Tohoku-Oki earthquake (11/03/2011–20/02/2021). The yellow star is the epicenter of the 2011 Tohoku-Oki earthquake. Dashed lines denote the trench axes of the Japan Trench and the Sagami Trough. Red and blue rectangular regions denote the downdip and updip regions, respectively, that are mentioned in the discussion. **c** Magnitude–time plot of the earthquakes (01/10/1997–20/02/2021) that occurred in the study region. The  $M_i \ge 4.0$  earthquakes are shown in green. The yellow star denotes the 2011 Tohoku-Oki earthquake

aftershocks occurring in the hanging wall owing to tensional stress changes and thrust aftershocks occurring in the footwall owing to compressional stress changes. Very few interplate earthquakes have occurred near the trench and in the outer rise, and the observed normal fault-type aftershocks were activated by tensional stress changes induced by the mainshock rupture (Asano et al. 2011; Obana et al. 2012, 2013).

Various physical models have been suggested to explain the spatiotemporal aftershock distribution of large earthquakes, with most models considering either coseismic stress changes (Dieterich 1994; King et al. 1994; Stein 1999; Toda and Stein 2003, 2022; Kroll et al. 2017; Ozawa and Ando 2021) or afterslip-induced stress loading (Marone et al. 1991; Perfettini and Avouac 2004, 2007; Perfettini et al. 2005; Hsu et al. 2006). However, the contribution of each mechanism to aftershock generation remains unclear. Here we investigate the aftershock generation mechanism following the 2011 Tohoku-Oki earthquake. We first apply the Hierarchical Space-Time Epidemic-Type Aftershock Sequence (HIST-ETAS) model (Ogata et al. 2003; Ogata 2004) to the decade of recorded seismicity since the 2011 Tohoku-Oki earthquake. Using the estimated HIST-ETAS model, we then separate the seismicity into background events (which are generated by either tectonic loading or aseismic phenomena) and triggered events (which are generated by earthquake-to-earthquake interactions). Finally, we discuss the plausible physical processes that generated the aftershocks following this megathrust event based on spatiotemporal variations in the background and triggered events.

## 2 Methods

We applied the HIST-ETAS model (e.g., Ogata et al. 2003; Ogata 2004; Bansal and Ogata 2013; Ueda et al. 2021) to the observed seismicity rate to investigate spatial variations in the seismicity characteristics and evaluate the probability of a given event being a background event. This model is a point process model that includes the Omori-Utsu law (Utsu 1961; Utsu et al. 1995), which formulates the typical aftershock temporal decay, the Utsu-Seki law (Utsu and Seki 1954), which formulates the relationship between the aftershock area and mainshock magnitude, the decay in aftershock density with distance (e.g., Felzer and Brodsky 2006), and a branching process, such that each earthquake, regardless of magnitude, has the ability to increase the probability of triggering a future earthquake (Iwata 2009).

The earthquake occurrence rate  $\lambda$  at time and location (t, x, y) and occurrence history  $H_t$  can be expressed by:

$$\lambda(t, x, y|H_t) = \mu + \sum_{t_i < t} \frac{K}{(t - t_i + c)^p} \left[ \frac{(x - x_i, y - y_i)S_i \left( \begin{array}{c} x - x_i \\ y - y_i \end{array} \right)}{\exp\{\alpha(M_i - M_c)\}} + d \right]^{-q},$$
(1)

where  $\mu$  is the background seismicity rate, and the second term expresses the rate of the earthquake occurrence triggered by a magnitude  $M_i$  earthquake at time and location  $(t_i, x_i, y_i)$ , where K, p, and c are the parameters of the Omori-Utsu law,  $S_i$  is a non-dimensional positive definite symmetric matrix for anisotropic clusters that is determined by identifying the aftershock cluster using a magnitude-based clustering algorithm (Ogata et al. 1995; Ogata 1998) and then choosing the best-fit ellipsoid that represents the cluster,  $\alpha$  is the aftershock magnitude sensitivity, q is the aftershock spatial decay rate, d is a constant, and  $M_c$  is the cutoff magnitude, 4.0 in this analysis. We note that  $H_t$  includes earthquakes that occurred before the target period (precursory period) because they are potentially influential to the seismicity in the target period (Ogata 2011).

The five seismicity parameters ( $\mu$ , K,  $\alpha$ , p, and q) are given as a function of space and are expressed as:

$$\mu(x, y) = \hat{\mu} \exp\{\vartheta_{\mu}(x, y)\}$$

$$K(x, y) = \hat{K} \exp\{\vartheta_{K}(x, y)\}$$

$$\alpha(x, y) = \hat{\alpha} \exp\{\vartheta_{\alpha}(x, y)\}, \qquad (2)$$

$$p(x, y) = \hat{p} \exp\{\vartheta_{p}(x, y)\}$$

$$q(x, y) = \hat{q} \exp\{\vartheta_{q}(x, y)\}$$

where  $\hat{\mu}$ ,  $\hat{K}$ ,  $\hat{\alpha}$ ,  $\hat{p}$ , and  $\hat{q}$  correspond to the geometric mean value of each parameter averaged over the analysis region. We adapt each  $\vartheta(x, y)$  value to the data by expressing each function using many coefficients that are placed in the locations of each earthquake epicenter and some additional points on the boundary of the analysis region. Each  $\vartheta(x, y)$  value at an arbitrary location is linearly interpolated using the three values at the vertices of each Delaunay triangle. The other parameters *c* and *d* are location independent.

The unknown parameters can be estimated via the maximum likelihood estimation method. The log-like-lihood is expressed as:

$$\log L = \sum_{k} \log \lambda (t_k, x_k, y_k | H_{t_k}) - \int_0^T \iint_S \lambda (t', x, y | H_{t'}) \mathrm{d}x \mathrm{d}y \mathrm{d}t',$$
(3)

where k is the index of each event in the analysis, S is the analysis region, and [0, T] is the analysis interval. However, it is hard to estimate the seismicity parameters stably because the number of unknown parameters is about five times larger than the number of data (events) used in the analysis. We obtained stable solutions by penalizing the spatial gradient of the parameter functions under the assumption of using smoothed functions (the facets of the piecewise linear function being as flat as possible):

$$Q(\theta|\tau) = \sum_{k=1}^{5} w_k \iint_A \left\{ \left(\frac{\partial \vartheta_k}{\partial x}\right)^2 + \left(\frac{\partial \vartheta_k}{\partial y}\right)^2 \right\} dx dy, \quad (4)$$

where  $\theta$  is a set of seismicity parameters, and  $\tau = (w_1, \ldots, w_5)$  is a set of weights that is used to tune the strength of the constraints. We then estimated the seismicity parameters that maximized the penalized log likelihood function (Good and Gaskins 1971):

$$R(\theta) = \ln L(\theta) - Q(\theta|\tau).$$
(5)

The weights  $\tau$  were objectively tuned by maximizing the integrated posterior distribution:

$$\Lambda(\tau) = \int_{\Theta} L(\theta) \pi(\theta | \tau) d\theta, \tag{6}$$

where  $\pi(\theta|\tau) \propto \exp\{-Q(\theta|\tau)\}$  is the probability distribution. The solution of  $\theta$  was then obtained for fixed weights  $\tau$  by maximizing the penalized log-likelihood in Eq. (5), which yields the optimal maximum a posteriori estimates. See Ogata et al. (2003) and Ogata (2004) for the details of the parameter estimation.

We estimated the uncertainties in the HIST-ETAS parameters at each location following the method outlined in Ueda et al. (2021). We first resampled the data 100 times by randomly extracting 90% of the earthquake data used in this analysis. Note that we did not exclude the M9 mainshock from the resampled data set. We then applied the HIST-ETAS model to each resampled dataset and estimated the spatial patterns of the five seismicity parameters. We note that the reference parameters ( $\hat{\mu}$ ,  $\hat{K}$ ,  $\hat{\alpha}$ ,  $\hat{p}$ , and  $\hat{q}$ ) may vary significantly among the 100 resampled datasets owing to the trade-offs between each parameter. Herak et al. (2001) investigated the interdependence of the ETAS model parameters using the aftershock sequences of the 1996 Ston-Slano earthquake and highlighted that the degree of correlation can be large, especially for pairs of aftershock parameters. Therefore, we normalized the model parameters by dividing the reference parameters estimated from each resampled dataset and then calculated the standard deviation of  $\vartheta(x, y) / \ln 10$  (common logarithm of the normalized parameter) at each location to evaluate the significance of the relative differences in each parameter, independent of the trade-offs between parameters.

We used the JMA unified hypocenter catalog (the preliminary Determination of Epicenters) as the earthquake data in our analysis. We applied the HIST-ETAS model to the  $M_j \ge 4.0$  earthquakes that occurred between the timing of the 2011 Tohoku-Oki mainshock (11 March 2011) and 20 February 2021 and were located at  $\le 100$  km depth beneath Japan (20–50° N, 120°–150° E). The  $M_j \ge 6.0$  earthquakes that occurred since 1922 and the  $M_j \ge 4.0$  earthquakes that occurred since February 2011 were used as the precursory occurrence history of the HIST-ETAS model.

## **3 Results**

The spatial distributions of the seismicity parameters that were estimated via the HIST-ETAS model, and their uncertainties are shown in Fig. 2. Here we focus on the spatial variations in the background seismicity rate  $\mu$  and purified aftershock productivity  $K/(c^p d^q)$ , which expresses the *i* -th event's aftershock intensity when and where the *i*-th event occurs in the Tohoku region (Fig. 2). The background seismicity rate after the 2011 Tohoku-Oki earthquake is low within the large coseismic slip zone (Yokota et al. 2011; Iinuma et al. 2012; Kato and Igarashi 2012), whereas it is high downdip of the large coseismic slip zone along the Pacific coast (Fig. 2a). These spatial differences are significant, given the low uncertainties (Fig. 2c). The purified aftershock productivity after the 2011 Tohoku-Oki earthquake is high around the large coseismic slip area (Fig. 2b). A comparison of the background seismicity rate and purified aftershock productivity suggests that triggered earthquakes have predominantly occurred near the large coseismic slip zone along the Japan Trench, where many aftershocks have occurred (Fig. 1). We don't discuss the other seismicity parameters (e.g.,  $\alpha$ ) because there are no significant spatial variations in these parameters in and around the large coseismic slip zone (Additional file 1: Fig. S1).

### 4 Discussion

## 4.1 Spatiotemporal distributions of background and triggered seismicity after the 2011 Tohoku-Oki earthquake

We find a significantly high background seismicity rate downdip of the large coseismic slip zone and relatively high aftershock productivity along the Japan Trench near the large coseismic slip zone compared with the background seismicity rate. Here we discuss the spatiotemporal distributions of the background and triggered seismicity after the 2011 Tohoku-Oki earthquake (Figs. 3, 4 and 5). We evaluate the numbers of background and triggered earthquakes in each  $0.1^{\circ} \times 0.1^{\circ}$  grid by calculating the summation of the probability that each event is classified into a background event and a triggered event, respectively, using the estimated HIST-ETAS parameters. The probability that each event is classified into a background event is explicitly given by Zhuang et al. (2002):

$$\varphi_i = \frac{\mu(x_i, y_i)}{\lambda(t_i, x_i, y_i | H_{t_i})},\tag{7}$$

(See figure on next page.)

**Fig. 2** Spatial distribution of the optimal maximum posterior estimates and their uncertainties of the HIST-ETAS model. **a** Common logarithm of the background seismicity rate  $\mu$ . **b** Common logarithm of the purified aftershock productivity  $K/(c^{p}d^{q})$ . **c** Standard deviations of the common logarithm of  $\mu$ . **d** Standard deviations of the common logarithm of  $K/(c^{p}d^{q})$ . Black dashed lines denote the Japan Trench axis. Blue solid lines denote the outer edge of the large coseismic slip zone of the 2011 Tohoku-Oki earthquake, as proposed by Kato and Igarashi (2012). Blue dashed and dotted lines denote the iso-slip contours (15-m contour) that were estimated by linuma et al. (2012) and Yokota et al. (2011), respectively



Fig. 2 (See legend on previous page.)



**Fig. 3** Spatiotemporal distributions of the number of background earthquakes calculated using the estimated HIST-ETAS model. **a**–**j** Spatial distributions of the number of background earthquakes by year. The black dashed line denotes the Japan Trench axis. The blue solid line denotes the outer edge of the large coseismic slip zone of the 2011 Tohoku-Oki earthquake, as proposed by Kato and Igarashi (2012)



Fig. 4 Spatiotemporal distributions of the number of triggered earthquakes calculated using the estimated HIST-ETAS model. **a**–**j** Spatial distributions of the number of triggered earthquakes by year. Black dashed and blue solid lines are the same as in Fig. 3



Fig. 5 Spatiotemporal distributions of the ratio of background seismicity to the total seismicity calculated using the estimated HIST-ETAS model. a-j Spatial distributions of the seismicity ratio by year. Black dashed and solid lines are the same as in Fig. 3

where the numerator and the denominator are the background seismicity rate and total seismicity rate when and where the i-th event occurs, respectively (Eq. 1). The probability that each event is classified into a triggered event is:

$$\rho_i = 1 - \varphi_i. \tag{8}$$

The spatiotemporal distributions of the background and triggered seismicity in the decade since the Tohoku-Oki earthquake (from 11 March 2011 to 11 March 2021) are shown in Figs. 3 and 4, respectively, and the spatiotemporal distribution of the ratio of background events to the total number of earthquakes is shown in Fig. 5. Along the Japan Trench near the large coseismic slip zone, few background earthquakes occurred during the analysis period (Fig. 3), whereas many triggered earthquakes occurred shortly after the 2011 Tohoku-Oki earthquake (Fig. 4a), followed by a decay in triggered seismicity over time (Fig. 4b-j); these spatiotemporal trends were expected based on the background seismicity rate and purified aftershock productivity results (Fig. 2a and b, respectively). The predominance of triggered seismicity relative to background seismicity is further confirmed in Fig. 5, thereby suggesting that the aftershocks along the Japan Trench after the 2011 Tohoku-Oki earthquake are largely due to earthquake-to-earthquake interactions. Many background and triggered earthquakes occurred downdip of the large coseismic slip zone right after the 2011 Tohoku-Oki earthquake (Figs. 3a and 4a), followed by decays in seismicity over time (Figs. 3 and 4). The background earthquakes contributed to the longterm evolution of seismicity following the Tohoku-Oki earthquake (Figs. 3b-j and 5), whereas the triggered seismicity temporarily increased when M7-class earthquakes (e.g., the 2016 and 2021 Fukushima-Oki earthquakes) occurred (Fig. 4f and j). We discuss the factors related to the temporal variations in background seismicity in the following section.

## 4.2 Relationship between background earthquakes and similar earthquakes downdip of the large coseismic slip zone

Previous studies have discussed the physical processes that control temporal variations in background seismicity (e.g., Hainzl et al. 2013; Kumazawa and Ogata 2013; Kumazawa et al. 2016; Nishikawa and Ide 2017; Nishikawa et al. 2021). Kumazawa and Ogata (2013) evaluated temporal variations in the background seismicity rate of the inland seismicity before and after the 2011 Tohoku-Oki earthquake using the non-stationary ETAS model and attributed the background seismicity rate changes to fluid-based aseismic triggering. Nishikawa and Ide (2017) and Nishikawa et al. (2021) detected earthquake swarms, which yielded sudden increases in the seismicity rate that could not be explained by the aftershock activity, and they suggested that aseismic phenomena, such as slow-slip events, could introduce temporal variations in the background seismicity rate. We compare the spatiotemporal variations in the background seismicity with those of similar earthquakes including repeating earthquakes, which indicate creep movement along the plate interface (Igarashi 2020), to discuss the physical processes that control the spatiotemporal variations in background seismicity downdip of the large coseismic slip zone. The spatial pattern of background seismicity resembles that of the similar earthquakes after the Tohoku-Oki earthquake (Fig. 6). The temporal variability in the number of background earthquakes downdip of the large coseismic slip zone (red rectangular region in Fig. 1) is shown in Fig. 7a, and a clear decay in background seismicity is observed. The relationship between the number of background earthquakes (Fig. 3) and aseismic slip amount, which was calculated using a similar earthquake catalog (Igarashi 2020), is shown in Fig. 7b. We used the method of Uchida et al. (2003) to calculate the aseismic slip amount from the similar earthquake catalog, whereby the slip amount of each similar earthquake is estimated from the following relationship between slip (d; cm)and the seismic moment  $(M_o; dyne \cdot cm)$  (Nadeau and Johnson 1998):

$$\log_{10} (d) = -2.36 + 0.17 \log_{10} (M_o).$$
<sup>(9)</sup>

The seismic moment for each event is estimated from its magnitude using the following relationship between magnitude and seismic moment (Hanks and Kanamori 1979):

$$\log_{10} \left( M_o \right) = 1.5M + 16.1. \tag{10}$$

The temporal variability in the number of background earthquakes closely follows a function of the form 1/t(Fig. 7a) and asymptotically approaches the background seismicity rate before the 2011 Tohoku-Oki earthquake, which is roughly estimated from Fig. 3 in Ogata (2011). Note that it will take more time to evaluate at what level the background seismicity will actually settle down to. The number of background earthquakes in each year positively correlates with the aseismic slip amount (Fig. 7b), which suggests that the increase in background seismicity downdip of the large coseismic slip zone is due to the afterslip of the 2011 Tohoku-Oki earthquake, which accumulates linearly with log time (Marone et al. 1991; Perfettini and Avouac 2007).



**Fig. 6** Spatial distributions of the numbers of background earthquakes and similar earthquakes during the 11/03/2011-11/03/2020 period. **a** Spatial distribution of the number of background earthquakes calculated using the estimated HIST-ETAS model. **b** Spatial distribution of the number of similar earthquakes ( $M_i \ge 4.0$ ) (Igarashi 2020). Black dashed and blue solid lines are the same as in Fig. 3

# 4.3 Implications for afterslip and aftershocks after the 2011 Tohoku-Oki earthquake

This study suggests that the spatiotemporal variations in the background events are due to the afterslip driven by the 2011 Tohoku-Oki earthquake (Fig. 7). The post-earthquake deformation following the Tohoku-Oki earthquake has been observed using land-based and seafloor geodetic instruments; however, the contributions of viscoelastic relaxation and afterslip to this deformation have been debated (e.g., Fukuda et al. 2013; Sun et al. 2014; Watanabe et al. 2014; Agata et al. 2019; Fukuda and Johnson 2021). Sun et al. (2014) indicated that numerical models of the transient viscoelastic mantle rheology could explain the landward motion of the trench area, thereby suggesting that viscoelastic relaxation plays the dominant role in short-term postseismic deformation, and Fukuda and Johnson (2021) separated the contributions of afterslip and viscoelastic relaxation under the assumption that these mechanisms are driven by coseismic stress changes and interact mechanically with each other to reveal the spatiotemporal afterslip distribution. Figure 7c shows a positive correlation between the number of background earthquakes (Fig. 3) and afterslip (Fukuda and Johnson 2021) averaged downdip of the large coseismic slip zone (red rectangle in Fig. 1). Although it is difficult to separate the contributions of afterslip and viscoelastic relaxation from the geodetic data, our study suggests that the contribution of afterslip can be extracted via a seismicity analysis of the aftershock distribution (Fig. 7).

Although many background and triggered earthquakes have occurred downdip of the large coseismic slip zone, triggered earthquakes were the predominant earthquake type right after the 2011 Tohoku-Oki earthquake (Figs. 3a, 4a, and 5a), thereby suggesting that these earthquakes were caused by the coseismic stress change. This result is consistent with the implication that the aftershocks in the surrounding regions were caused by the stress concentration due to the large coseismic slip (Asano et al. 2011; Kato and Igarashi 2012). Conversely, background earthquakes are the predominant earthquake type in the long term (Figs. 3b-j and 5b-j), and they positively correlate with the aseismic slip amount (Fig. 7b and c), thereby suggesting that temporal variations in the tectonic loading rate (afterslip) have a greater impact on the physical aftershock generation mechanism than the coseismic stress change due to the mainshock



**Fig. 7** Temporal variations in the number of background earthquakes compared with the aseismic slip amount. **a** Temporal variations in the number of background earthquakes downdip of the large coseismic slip zone (red rectangle in Fig. 1). The blue curve is the best fit of the function a + b/(t + c) to the observed number of background earthquakes per year (red points). The two black dashed lines provide a possible range for the number of background earthquakes per year before the 2011 Tohoku-Oki earthquake, which is deduced from Fig. 3 in Ogata (2011). **b** Relationship between the number of background earthquakes calculated using the estimated HIST-ETAS model and the aseismic slip amount, which was calculated using the similar earthquake catalog (Igarashi 2020) and the method outlined by Uchida et al. (2003). Black dashed lines are the same as in **a**. **c** Relationship between the number of background earthquakes calculated using the estimated HIST-ETAS model and the afterslip, which was estimated via a geodetic inversion analysis (Fukuda and Johnson 2021). Black dashed lines are the same as in **a** 

rupture (Fig. 8). Predominance of triggered earthquakes right after the 2011 Tohoku-Oki earthquake and their faster decay compared with background earthquakes can also be confirmed by temporal variations in background and triggered earthquakes (Additional file 1: Fig. S2).

However, the reverse trend is observed updip of the large coseismic slip zone, as triggered earthquakes are the predominant earthquake type in the long term (Figs. 3, 4, and 5), thereby suggesting that the aftershocks are caused by earthquake-to-earthquake interactions,

and that there is a lack of afterslip in this updip region along the Japan Trench (Fig. 8). This result supports previous findings of the normal fault aftershocks in the outer rise potentially being activated by an abrupt increase in extensional stress caused by the mainshock rupture (earthquake-to-earthquake interactions, Asano et al. 2011; Obana et al. 2012, 2013), the lack of afterslip in the updip region (Sun et al. 2014; Agata et al. 2019), and the observation of very few interplate earthquakes in the updip region (Obana et al. 2013, 2021).



Fig. 8 Interpretation of the aftershock generation mechanisms following the 2011 Tohoku-Oki earthquake. The red line denotes the plate interface (Nakajima and Hasegawa 2006; Kita et al. 2010), and the gray circles are the hypocenters of the events in the F-net hypocenter catalog (11/03/2011–11/03/2021)

#### 4.4 Previous aftershock models

Various physical aftershock models have been proposed to explain aftershock activity, which can be divided into two main categories: static stress changes induced by the mainshock and large aftershocks (Dieterich 1994; King et al. 1994; Stein 1999; Toda and Stein 2003, 2022; Kroll et al. 2017; Ozawa and Ando 2021) and afterslip-induced stress loading (Marone et al. 1991; Perfettini and Avouac 2004, 2007; Perfettini et al. 2005; Hsu et al. 2006). Toda and Stein (2022) forecasted the seismicity changes after the Tohoku-Oki earthquake based on a rate-and-state friction model (Dieterich 1994) that was coupled with Coulomb stress transfer, and they qualitatively explained the quiescence in the mainshock source region and activation in the surrounding area, whereas Perfettini et al. (2019) described the aftershock migration away from the rupture area of the Tohoku-Oki mainshock using an analytical model that was based on afterslip-induced seismicity. Although these studies only considered either mainshock-induced stress changes or afterslip-induced stress loading, our results suggest the importance of considering both mechanisms to explain the aftershock sequence of the 2011 Tohoku-Oki earthquake. Many aftershocks have occurred updip of the large coseismic slip zone along the Japan Trench (blue rectangular region in Fig. 1) without any detectable afterslip (Sun et al. 2014; Agata et al. 2019), whereas the aftershock activity downdip of the large coseismic slip zone (red rectangular region in Fig. 1) has been affected by both mainshockinduced Coulomb stress transfer and afterslip-induced stress loading ((Figs. 3, 4, 5, and 7). Our study successfully divides the aftershocks into background earthquakes, which are inferred to be afterslip-induced events (Fig. 3), and triggered earthquakes, which are inferred to be induced by the stress transfer from other earthquakes (Fig. 4). The present classification is consistent with both the afterslip distribution (Fukuda and Johnson 2021) and mainshock-induced stress change (Asano et al. 2011). Our results suggest that the Coulomb stress model and afterslip model should be combined in future studies to better understand aftershock activity. Furthermore, applying our method to the aftershock sequences of other large earthquakes and quantifying the afterslip-induced seismicity may provide new insights into the relationship between afterslip and aftershocks (Churchill et al. 2022). Our method can be also applied to estimate the afterslip evolution if earthquake catalogs are available, even in areas where geodetic observation networks are sparse. It is therefore possible to employ the classifications using the HIST-ETAS model to advance our understanding of the seismicity characteristics following large earthquakes and to forecast the spatiotemporal evolution of both background and triggered earthquakes. Furthermore,

## **5** Conclusions

We investigated the spatiotemporal characteristics of the aftershocks following the 2011 Tohoku-Oki earthquake by applying the HIST-ETAS model to the decade of recorded seismicity since the mainshock. We successfully divided aftershocks into background earthquakes, which are afterslip-induced events, and triggered events, which are induced by the stress transfer from other earthquakes. Triggered earthquakes are the predominant earthquake type updip of the large coseismic slip zone along the Japan Trench, which is consistent with the lack of afterslip in this region. Background earthquakes are the predominant earthquake type downdip of the large coseismic slip zone, and they positively correlate with the afterslip evolution that was detected using repeating earthquakes and geodetic data. Our results suggest that both stress transfer and afterslip should be considered for both modeling an aftershock sequence and evaluating the relationship between afterslip and aftershocks.

#### Abbreviations

HIST-ETAS Hierarchical Space–Time Epidemic-Type Aftershock Sequence JMA Japan Meteorological Agency

#### Supplementary Information

The online version contains supplementary material available at https://doi. org/10.1186/s40645-023-00564-0.

Additional file 1: Figure S1. Spatial distribution of the optimal maximum posterior estimates and their uncertainties of the HIST-ETAS model.Common logarithm of *a*. Common logarithm of *p*.Standard deviations of the common logarithm of *a*. Standard deviations of the common logarithm of *p*. Black dashed lines denote the Japan Trench axis. Blue solid lines denote the Japan Trench axis. Blue solid lines denote the outer edge of the large coseismic slip zone of the 2011 Tohoku-Oki earthquake, as proposed by Kato and Igarashi. Blue dashed and dotted lines denote the iso-slip contoursthat were estimated by linuma et al.and Yokota et al., respectively. Figure S2. Temporal variations in the number of background and triggered events. Blue circles show temporal variations in the number of triggered events in the red rectangular region in Fig. 1. Red circles show temporal variations in the number of background events in

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### Author contributions

TU applied the HIST-ETAS model to the earthquake catalog. TU and AK interpreted the results. TU drafted the manuscript. All of the authors read and approved the final manuscript.

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#### Availability of data and materials

The JMA and F-net catalogs are available from the NIED (National Research Institute for Earth Science and Disaster Resilience) Data Management Center (https://hinetwww11.bousai.go.jp/auth/?LANG=en and https://www.fnet. bosai.go.jp/event/search.php?LANG=en, respectively). The source codes (FORTRAN) and manual for the HIST-ETAS model can be found in Ogata et al. (2021).

## Declarations

#### **Competing interests**

The authors declare that they have no competing interests.

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